

Appendix: Quantitative modeling of the molecular steps underlying shut-off of rhodopsin's activity in rod phototransduction

Trevor D. Lamb¹ and Timothy W. Kraft²

¹Eccles Institute of Neuroscience, John Curtin School of Medical Research, The Australian National University, Canberra, ACT 2601, Australia

²Department of Optometry and Vision Science, University of Alabama at Birmingham, Birmingham, AL 35294, USA

A Boundary condition at photoisomerization position

For free aqueous diffusion of cGMP, Fick's Second Law is

$$\frac{\partial cG}{\partial t} = D \frac{\partial^2 cG}{\partial x^2}. \quad (\text{A.1})$$

Within the rod outer segment, though, longitudinal diffusion is 'baffled' by the stack of disks, and as a result only a fraction f_A of the cross-sectional area A is available for diffusion, and only a fraction f_V of the envelope volume is cytoplasmic [1]. Furthermore, as invoked by Lamb & Pugh [2], the cytoplasm is envisaged to have a buffering capacity BP for cGMP. As a result, the modified diffusion equation is

$$\frac{\partial cG}{\partial t} = D_{cG} \frac{\partial^2 cG}{\partial x^2} \quad (\text{A.2})$$

where the effective longitudinal diffusion coefficient D_{cG} is

$$D_{cG} = \frac{f_A}{f_V BP} D. \quad (\text{A.3})$$

From Fick's First Law, the longitudinal flux of cGMP along the outer segment (which can occur only via the available area $f_A A$) is given by

$$\text{Longitudinal flux} = -D f_A A N_{Av} \frac{\partial cG}{\partial x} \quad (\text{A.4})$$

where the true diffusion coefficient D , rather than the effective diffusion coefficient D_{cG} , must be used, and where Avogadro's number N_{Av} converts the units to molecules s^{-1} . At a disk that has received a single photoisomerization, the rate of hydrolysis of cGMP is

$$\text{Hydrolysis rate} = E^*(t) \frac{k_{cat}}{K_m} cG(x, t) = E^*(t) \frac{\frac{1}{2}k_{CAT}}{K_m} cG(x, t) \quad (\text{A.5})$$

where $k_{cat} = \frac{1}{2}k_{CAT}$ is the maximal rate of hydrolysis by a single $G^* - E^*$, and k_{CAT} is the maximal rate of hydrolysis by the fully-activated PDE dimer.

In the symmetrical case, with a single photoisomerization occurring at the middle of the outer segment ($x = x_0$), the magnitude of the unidirectional flux on either side of this point is half the total hydrolytic rate, so that from Eqns (A.4) and (A.5)

$$D f_A A N_{Av} \left. \frac{\partial cG}{\partial x} \right|_{x_0^+} = \frac{1}{2} E^*(t) \frac{\frac{1}{2}k_{CAT}}{K_m} cG(x_0, t). \quad (\text{A.6})$$

From Eqn (A.3) we have

$$D f_A = D_{cG} f_V B P D \quad (\text{A.7})$$

and substitution into Eqn (A.6) gives

$$D_{cG} f_V A B P N_{Av} \left. \frac{\partial cG}{\partial x} \right|_{x_0^+} = \frac{1}{2} E^*(t) \frac{\frac{1}{2}k_{CAT}}{K_m} cG(x_0, t) \quad (\text{A.8})$$

or

$$\left. \frac{\partial cG}{\partial x} \right|_{x_0^+} = \frac{E^*(t)}{2 D_{cG}} \left[\frac{\frac{\frac{1}{2}k_{CAT}}{K_m}}{f_V A B P N_{Av}} \right] cG(x_0, t). \quad (\text{A.9})$$

Formulation of Lamb & Pugh (1992)

For the term in [] above, we note that Lamb & Pugh [2] defined β_{sub} in their Eqn (4.4) as

$$\beta_{sub} = \frac{\frac{\frac{1}{2}k_{CAT}}{K_m}}{V_{cyto} N_{Av} B P} \quad (\text{A.10})$$

where the cytoplasmic volume V_{cyto} is given by

$$V_{cyto} = f_V A L. \quad (\text{A.11})$$

Hence

$$\beta_{sub} L = \frac{\frac{\frac{1}{2}k_{CAT}}{K_m}}{f_V A N_{Av} B P} \quad (\text{A.12})$$

where the right hand side of the equation above is the term in [] in Eqn (A.9). Hence the latter equation may be rewritten as

$$\left. \frac{\partial cG}{\partial x} \right|_{x_0^+} = \frac{E^*(t) \beta_{sub} L}{2 D_{cG}} cG(x_0, t) \quad (\text{A.13})$$

which is presented as Eqn (4.11) in the Theory section, and which is exactly equivalent to Eqn (B 1) of Lamb & Pugh [2].

Formulation of Gross, Pugh & Burns (2012)

In a variant of the symbols, Gross et al [3] defined β_{idv} (see p. 1781) as

$$\beta_{\text{idv}} = \frac{\frac{k_{\text{cat}}}{K_m}}{V_{\text{id}} N_{\text{Av}}} = \frac{\frac{\frac{1}{2}k_{\text{CAT}}}{K_m}}{V_{\text{id}} N_{\text{Av}}} \quad (\text{A.14})$$

where the interdiscal cytoplasmic volume V_{id} is given by

$$V_{\text{id}} = f_V A \delta \quad (\text{A.15})$$

with $\delta = L/N_{\text{disks}}$ being the mean interdisk spacing. Hence

$$\beta_{\text{idv}} \delta = \frac{\frac{\frac{1}{2}k_{\text{CAT}}}{K_m}}{f_V A N_{\text{Av}}} \quad (\text{A.16})$$

or

$$\frac{\beta_{\text{idv}} \delta}{BP} = \frac{\frac{\frac{1}{2}k_{\text{CAT}}}{K_m}}{f_V A BP N_{\text{Av}}} \quad (\text{A.17})$$

where the right hand side is again the term in [] in Eqn (A.9). Accordingly, Eqn (A.9) may be rewritten in the terminology of Gross et al [3] as

$$\left. \frac{\partial cG}{\partial x} \right|_{x_0^+} = \frac{E^*(t) \beta_{\text{idv}} \delta}{2 D_{\text{cG}} BP} cG(x_0, t). \quad (\text{A.18})$$

We noticed that this equation differs from Eqn (6) of [3], which instead gives the denominator as $4 D_{\text{cG}}$. As the buffering power BP is not mentioned anywhere in Gross et al [3], it is clear that they assumed $BP = 1$. Accordingly, in their terminology, the required boundary condition is

$$\left. \frac{\partial cG}{\partial x} \right|_{x_0^+} = \frac{E^*(t) \beta_{\text{idv}} \delta}{2 D_{\text{cG}}} cG(x_0, t). \quad (\text{A.19})$$

Eqn (6) of Gross et al [3] differs from this equation by a factor of 2, and in our view this occurred because they incorrectly invoked an additional volume fraction of $\frac{1}{2}$. As a result, their subsequent analysis of ‘rogue’ responses actually calculated $2\beta_{\text{idv}}$, instead of β_{idv} .

Finally, we note that

$$\beta_{\text{idv}} \delta = \beta_{\text{sub}} L \quad (\text{A.20})$$

so that

$$\beta_{\text{idv}} = N_{\text{disks}} \beta_{\text{sub}} \quad (\text{A.21})$$

where $N_{\text{disks}} = L/\delta$ is the number of disks in the outer segment.

Appendix References

1. Lamb TD, McNaughton PA, Yau K-W. Spatial spread of activation and background desensitization in rod outer segments. *J Physiol* 1981; 319:463-496.
2. Lamb TD, Pugh EN Jr. A quantitative account of the activation steps involved in phototransduction in amphibian photoreceptors. *J Physiol* 1992; 449:719-757.
3. Gross OP, Pugh EN Jr, Burns ME. Spatiotemporal cGMP dynamics in living mouse rods. *Biophys J* 2012; 102:1775-1784.